

Do problems 5.0 and 5.1.

Problem 5.0: Consider the complex scalar field of Eq.(5.2.11). Compute the mean value in the vacuum of the time-ordered product

$$-i \Delta(x, y) = (\Phi_0, T \{ \phi(x) \phi^\dagger(y) \} \Phi_0) = \langle 0 | T \{ \phi(x) \phi^\dagger(y) \} | 0 \rangle \quad (1)$$

or

$$-i \Delta(x, y) = \theta(x^0 - y^0) \langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle + \theta(y^0 - x^0) \langle 0 | \phi(y^\dagger) \phi(x) | 0 \rangle \quad (2)$$

in which all quantities are those of free fields. You may use Eq.(6.2.15) to express $\theta(x^0 - y^0)$ as an integral over an energy variable s and then change variables to $q = (\vec{p}, p^0 + s)$ where $p^0 = \sqrt{\vec{p}^2 + m^2}$. You may then use the identity

$$\theta(-t) = \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{e^{-ist}}{s - i\epsilon} ds \quad (3)$$

to write $\theta(y^0 - x^0)$ as an integral over an energy variable s and then change variables to $r = (-\vec{p}, s - p^0)$. In both cases it will behoove you to imitate Weinberg's arguments on pages 276 and 277.

Hints for problem 5.1: You may need to use Eq.(2.5.10). Problem 5.1 is easier in matrix notation. If we suppress the species index n , which plays no role, and introduce the matrix

$$U_{i,\sigma}(\vec{p}) = u_i(\vec{p}, \sigma), \quad (4)$$

then we may write Eq.(5.1.23) as

$$U(\vec{0}) D^{(j)}(R) = D(R) U(\vec{0}). \quad (5)$$

Eq.(5.1.21) then is

$$U(\vec{q}) = \left(\frac{m}{q^0}\right)^{1/2} D(L(q)) U(\vec{0}), \quad (6)$$

and Eq.(5.1.19) is

$$U(\vec{p}_\Lambda) D^{(j)}(W(\Lambda, p)) = \left(\frac{p^0}{(\Lambda p)^0}\right)^{1/2} D(\Lambda) U(\vec{p}). \quad (7)$$

In these formulae the matrices D without the superscript (j) are (in general, non-unitary) matrices in a representation of the Lorentz group. The matrices $D^{(j)}(R)$ are the usual unitary matrices that represent the (compact) rotation group.